IMPLICATIONS OF ELECTRON BEAM PROPAGATION ON **SOLAR FLARE ENERGETICS**

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Standard and extended models of solar eruptive events

- Standard model with return currents and open questions
- How do return currents affect HXR observations?
- Motivation for developing RUNAWAYRC

Recent advances in modeling the beam/return current propagation

- Beam/return current model which neglects runaway electrons Allred+20 (FP) and which includes runaway electrons Alaoui+21 (RUNAWAYRC)
- Thermal response to the injection of an electron beam including return currents

Outline

STANDARD MODEL FOR SOLAR ERUPTIVE EVENTS

Cartoon of 2D standard model



Real observations

"PROBLEMS" WITH STANDARD MODEL

IIMBER "PROBLEM"

Electron flux required to produce observed X-ray emission is ~10³⁶ electrons/s. (Hoyng et al 1976, 1978)

Number of electrons in typical flaring region $nV \sim 10^{37}$ => evacuation of the flaring region in seconds

CURRENT STABILITY

Currents 10^{17} Amps Induced magnetic field 10⁵ G => pinched beam (Bennet 1934, 1955)

Fast electron beam





Extended Standard model for electron propagation with return currents



Co-spatial Return Current

Which electrons carry the return current?



- (1) Thermal
- (2) Suprathermal Runaway Electrons
- (3) Nonthermal beam electrons scattered backwards



Outstanding questions

Is there a feedback between the large scale propagation effects and the acceleration process itself?

How much and where is the electrons energy dissipated?

Do electron beams accelerated in the solar corona produce the observed spectral properties?



How do return current losses affect X-ray spectra?



Lowest energy electrons lose the highest *fraction* of their energy

Flatten electron distribution at low-energies

A large enough value of the low energy cutoff is also observed as a flattening in the X-ray spectrum

Question

To what extent is the shape of X-ray spectra a consequence of the acceleration mechanism or beam propagation?



Mechanisms which affect the HXR emission



With imaging spectroscopy with higher range of sensitivity (dynamic range), high spatial and temporal resolution: can differentiate between these effects

Are current-driven instabilities responsible for X-ray flattening?



Classical (Spitzer) resistivity proportional to $T^{-3/2}$

Resistivity up to 7 orders of magnitude higher than Spitzer



Alaoui & Holman 2017





 v_{Te} min) (Е_с >0

RC driven instabilities can occur in solar flare conditions but they are not sufficient to explain the flattening at lower energies in X-ray spectra in most cases



Observational and theoretical motivation for runaway model

Electric field strength from models without runaways



Kinetic modeling of electron beam propagation in flares





Cold target: Energy of the beam electrons>> energy of plasma with which they interact: It does not mean that the plasma has a low T

	Applicability con	References		
	High value of low energy cutoff	Hot	Brown 71, 73; Emslie 78	
et)	Any low energy cutoff	plasma	Jeffrey+2014, 2019; Kontar+ 2015, Emsl	
	High value of low energy cutoff	Any	Zharkova et al 95, 10; Zharkova & Gordov Siversky & Zharkova 2009 Knight & Stu	
	Any low energy	temperature	Holman 2012 Allred, Alaoui, Kowalski, Kerr 2020	
	cuton		Alaoui, Holman, Allred, Eufrasio 2021	
yay		Any plasma co-spatial	Not modeled self-consistent	



Model Description: Alaoui et al. 2021

POWER-LAW INJECTED

accelerated electrons continuously injected at apex of 1D loop model

CURRENT BALANCE

 $J_{beam}(x) = J_{RC}(x)$ Ohm's law

STEADY-STATE

Time scales >> than electron-ion collision time, i.e., return current/beam system reached steady-state (Van den Oord 1990, Siversky & Zharkova 2009)

BOUNDARY CONDITION

STABLE RETURN CURRENT

SUB-DREICER ELECTRIC FIELD

THERMALIZATION OF ELECTRONS

No current-driven instabilities. Resistivity is Spitzer

We use higher values for E_{RC} but the accuracy of the solution decreases with increasing E_{RC}

If energy of direct beam electrons reaches thermal energy, electrons lost from beam

Runaway growth rate from Landreman et al. (2014) $\int J_{drift}(x) + J_{runaway}(x)$

No runaway electrons at the footpoints. $J_{runaway}(x = L) = 0$

 $E_{RC} < 0.11 E_D$ everywhere along the electrons' path





Return current affects acceleration region and chromosphere



Current density linearly proportional to flux density



Main implications

- (1) 43% of flux returning to acceleration region is suprathermal (energy gain 21 keV)
- (2) Electron flux injected into chromosphere reduced due to thermalization by the return current



Electric field & potential drop spatial evolution



Energy of runaway electrons at looptop (gain of 21 keV) >> thermal energy Electrons returning to the acceleration region are already suprathermal=> further accelerated to keep acceleration ongoing





Return current energy losses dominate over Coulomb collisions



Return currents cannot be neglected

Even when considering heating reduction due to presence of runaways



Lower low-energy cutoff results in thermalization of more beam electrons in the corona=>reduced electron flux into chromosphere

In hotter plasmas return currents are still significant

Same atmosphere with apex temperature 10 MK, same spectral index $\delta=4$



RC significant but runaways negligible (5% at LT)

to example on left

Higher injected flux=> higher runaways and higher reduction of heating +Coulomb collisions contribute to reducing the heating especially in runaway case

Same injected flux density compared to example in the middle, higher low-energy cutoff

Beam electrons thermalized below transition region





Fraction of runaways at the looptop: Various regimes for propagation

Apex density
$$n_e = 8 \times 10^9 \ cm^{-3}$$

$$\delta = 4$$
; $E_c = 20$ keV



Under which conditions are the return current losses negligible? Under which conditions are runaways significant?



Fraction of injected energy flux reaching the transition region is reduced by the return current



How much of the injected energy flux density at the corona reaches the top of the transition region?

Initial apex temperature Top panels: 3 MK Lower panels: 10 MK

Misinterpretation of the energy flux injected into the transition region and therefore also the accelerated flux density injected at the top of the loop



How neglecting return current losses affects the thermal response





The heating in the chromosphere is overestimated when return currents are not considered





Many large flares are associated with higher flux densities than previous examples

Recent modeling which requires accounting for the return current energy losses

Range of injected fluxes between vertical red lines Kowalski et al 2022

model ID	Beam Flux [erg s^{-1} cm ⁻²]	δ	$E_c \; [\text{keV}]$	μ_o	Beam Duration [s]	Comment
c15s-5F11-25-4.2	5×10^{11} (5F11)	4.2	25	0.1	15	"Extended heating" 5F11 in Paper I.
c20s-F11-25-4	10^{11} (F11)	4	25	0.1	20	Published in Kuridze et al. (2020) .
c20s-F11-15-5	10^{11} (F11)	5	15	0.1	20	Published in Graham et al. (2020) .

Only lower energy cutoffs are consistent with observations Graham et al 2020 BUT the chromosphere gets heated faster in the simulations

Because the return current results in a significant energy deposition in the corona, can timing inconsistencies between models and observations be resolved simply by including them in the simulations?

What constraints on the beam energy distribution and the initial atmosphere can be deduced?

Do electron beams accelerated in the solar corona produce the observed spectral properties?







Understanding electron beam/return current propagation is important for: **THERMAL RESPONSE**

(1) Corona is heated faster and to higher temperatures; below TR stays cooler for longer (2) The injected energy (and flux) at the looptop is significantly different from that injected into the chromosphere

Flares with high injected flux densities should be reanalyzed to include the return current effects

ELECTRON BEAM ACCELERATION

Runaways provide suprathermal particles to the looptop. Are particles accelerated there? If so, the runaways are seed particles for continuing acceleration

Runaways energies ~10-50 keV; runaway fractions can be tens of % at the looptop



