

GRBs from NS+BH coalescences: expected rates and properties



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Coalescing compact binaries

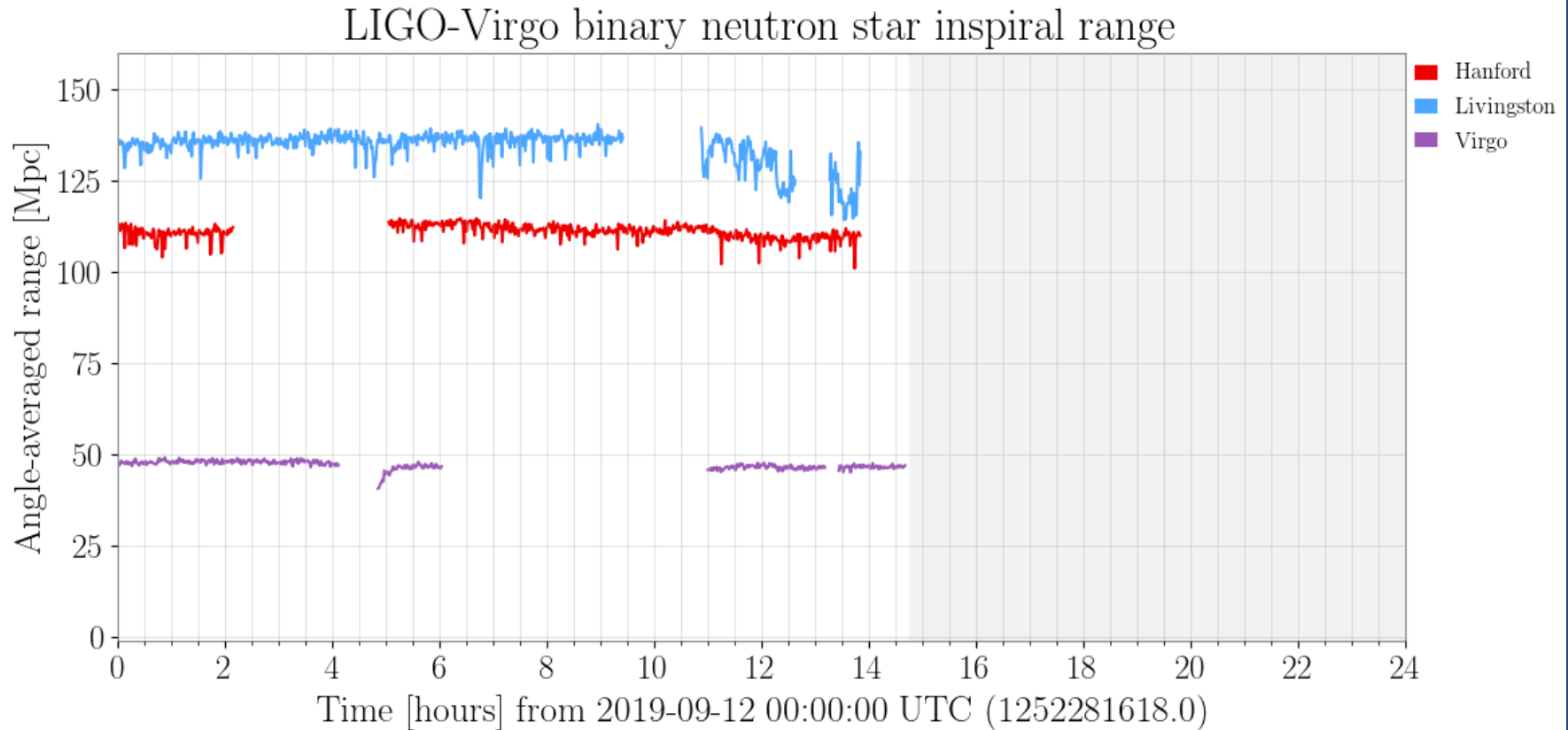
- Waveform from two coalescing point-like masses is determined by a combination of component masses (the chirp mass)

$$M_{ch} = (\mu^3 M^2)^{1/5}$$

$$h \sim M_{ch}^{5/3} f^{2/3} / r$$

$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}} = \frac{c^3}{G} \left(\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right)^{3/5}$$

Current Detection horizon



$$D_h \sim M_{\text{chirp}}^{5/6}$$

<https://www.gw-openscience.org/>

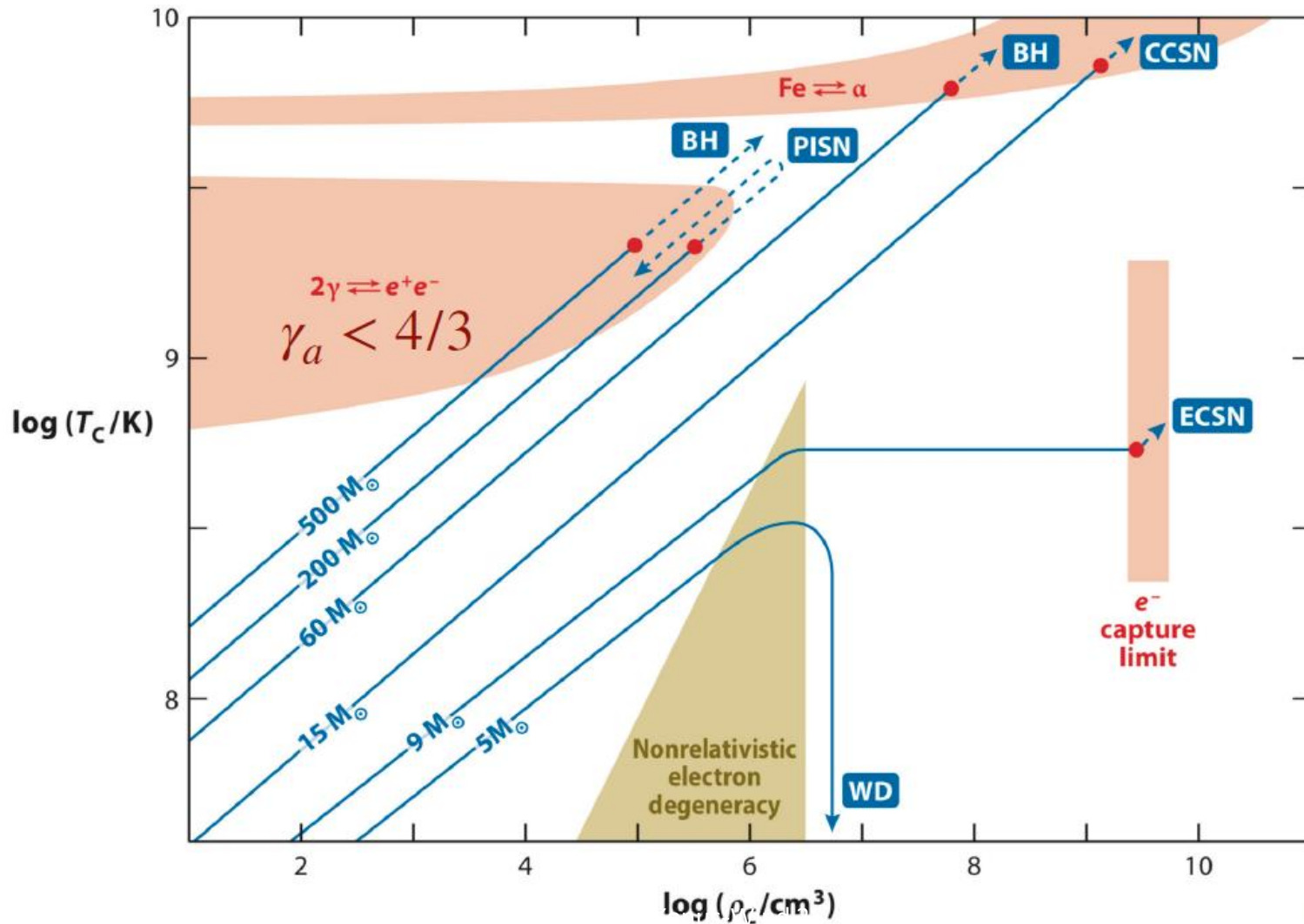
LVC O3 detections

- 25 triggers, 20 BH+BH, 3 NS+NS, 2 NS+BH
<https://gracedb.ligo.org>
- No electromagnetic counterparts so far

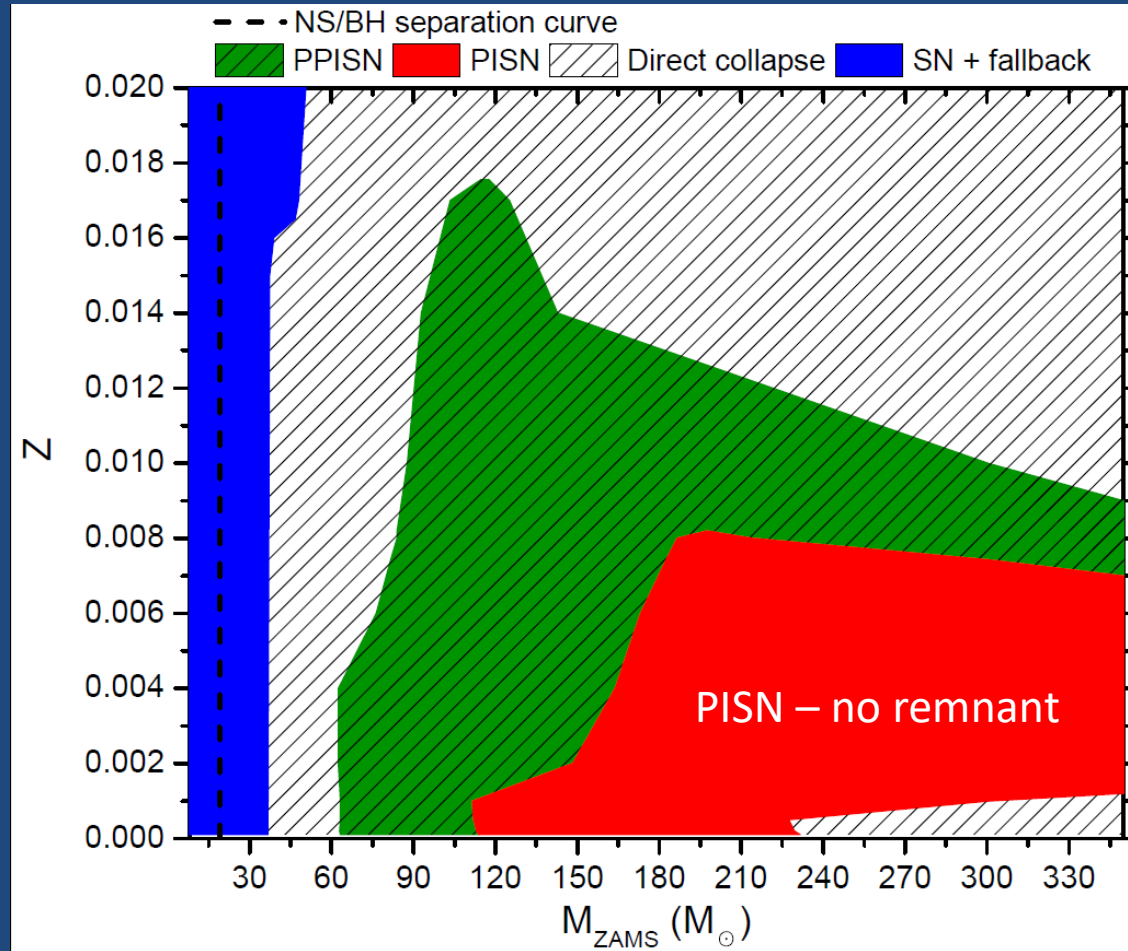
Astrophysical binary BHs

- BH formation
 - Initial ZAMS mass
 - Initial rotation
 - Possible kick
- BH in binaries
 - From massive binary systems
 - In dense stellar clusters (dynamical)
 - Primordial BHs

BH formation from stars (solar metallicity)



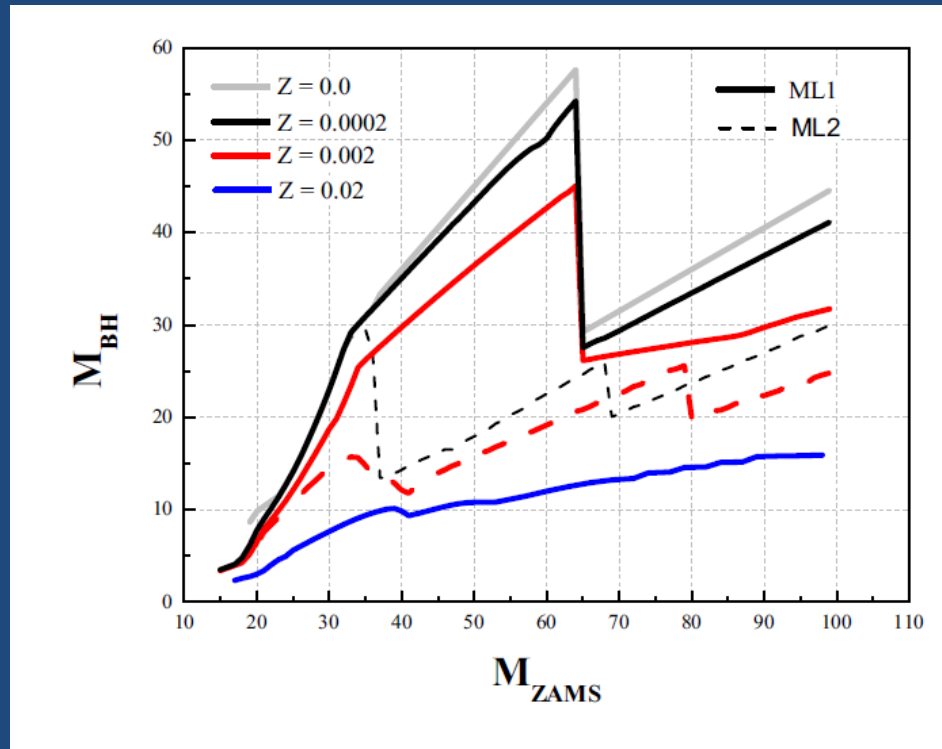
BH from massive stars



Spera, Mapelli
2017

Pair instability SNe (PISNe, Fowler & Hoyle 1964), pulsational PISNe (PPISNe) (Woosley 2017)

Dependence on the metallicity and stellar wind mass loss



Fryer et al 2012, Giacobbo et al. 2018

Additional effects in binaries:

- Initial spin misalignment
- Tidal synchronization of the envelope
- Common envelope phase
- Star formation and metallicity history in galaxies

$$\Psi \left(z, \frac{Z}{Z_{\odot}} \right) = \psi(z) \Phi(Z/Z_{\odot}).$$

$$\Phi(Z/Z_{\odot}) = \frac{\hat{\Gamma}[0.84, (Z/Z_{\odot})^2 10^{0.3z}]}{\Gamma(0.84)}$$

M1		M2	A
50.00		36.00	190.00
48.49		34.25	197.50
46.03		34.09	203.90
28.50		47.50	235.60
28.50		47.50	235.60
23.91	WR	52.09	278.40
	SN Ib		
10.76	BH	52.09	347.80
10.76	BH	51.91	348.80
10.76	BH	49.26	364.30
10.76	BH	44.88	208.30
10.76	BH	25.32	12.38
			WR
	BH		SN Ib
10.76	BH	11.40	27.60
			BH
			Coalescence
		22.16	BH

- Initial ZAMS mass
- Fraction of collapsed mass

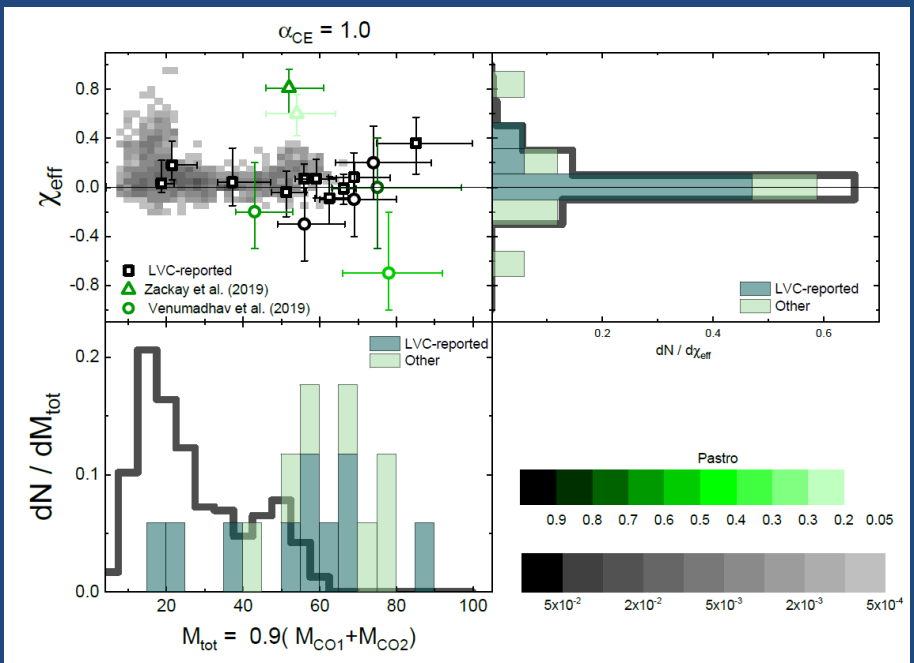
$$k_{BH} = M_{BH}/M_*$$

- Possible kick

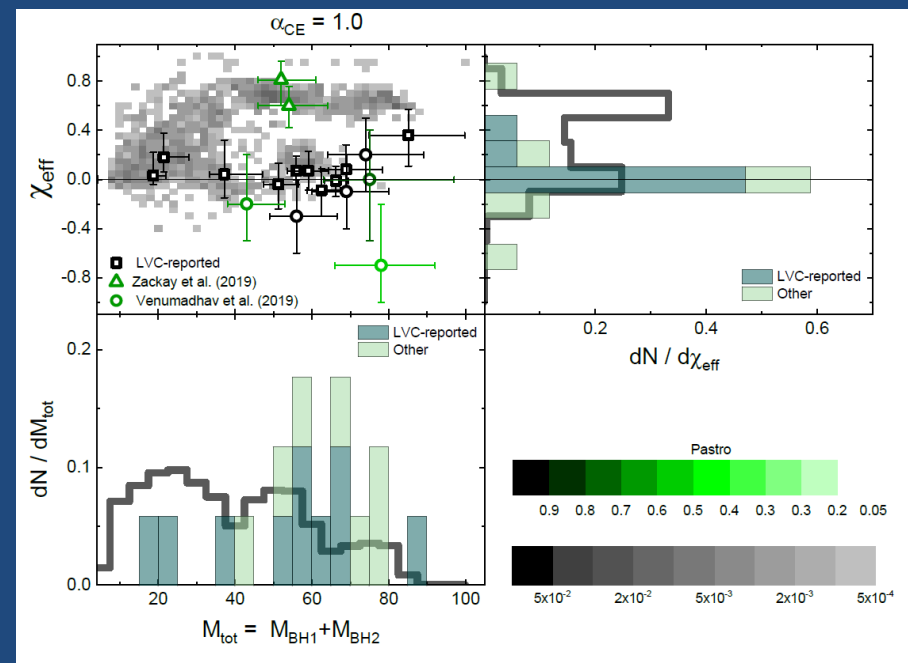
$$\frac{w_{BH}}{w_{NS}} = \frac{M_* - M_{BH}}{M_* - M_{OV}} = \frac{1 - k_{BH}}{1 - M_{OV}/M_*}$$

Tutukov, Yungelson 1993 MNRAS)
 Lipunov, Postnov, Prokhorov 1997 MNRAS)

Mass and spin distributions before coalescence



Without fallback

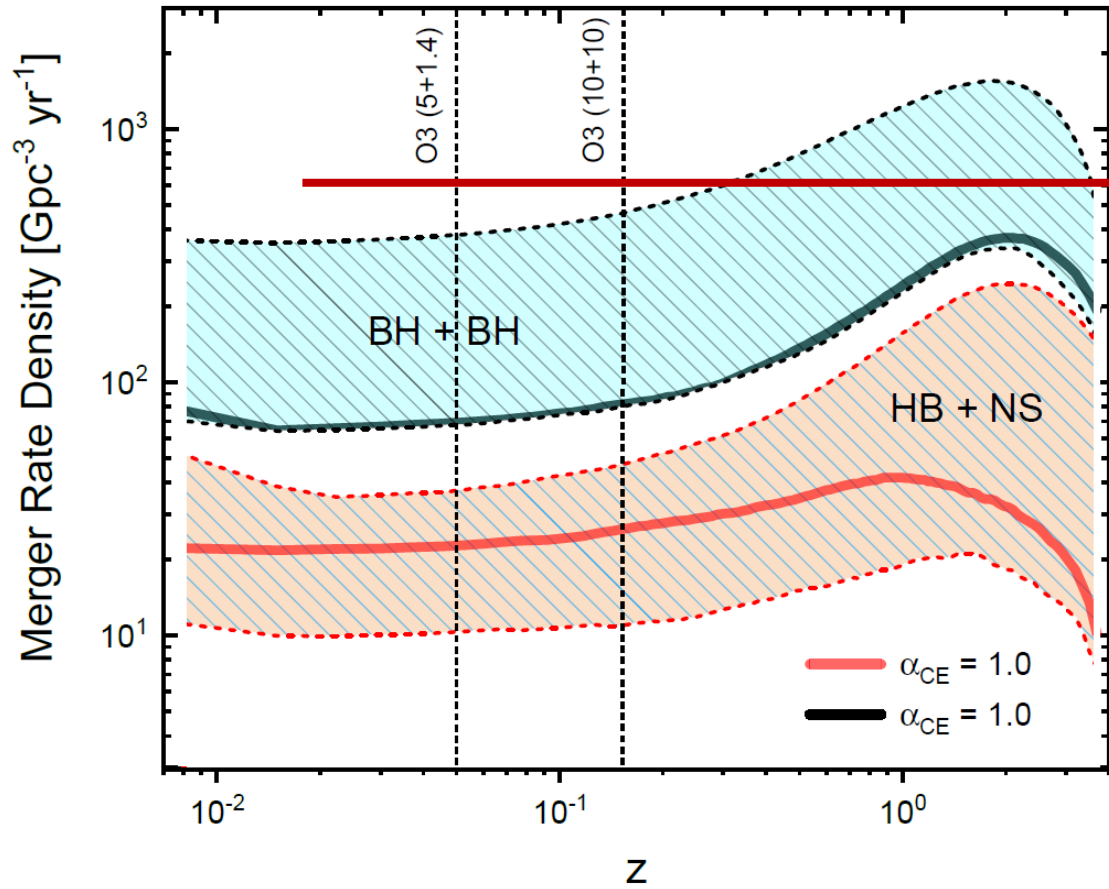


+ fallback from envelope

PK+ 2019, Physics-Uspekhi (2019, No 11, in press)

MNRAS 483, 3288–3306 (2019)

BH+NS systems

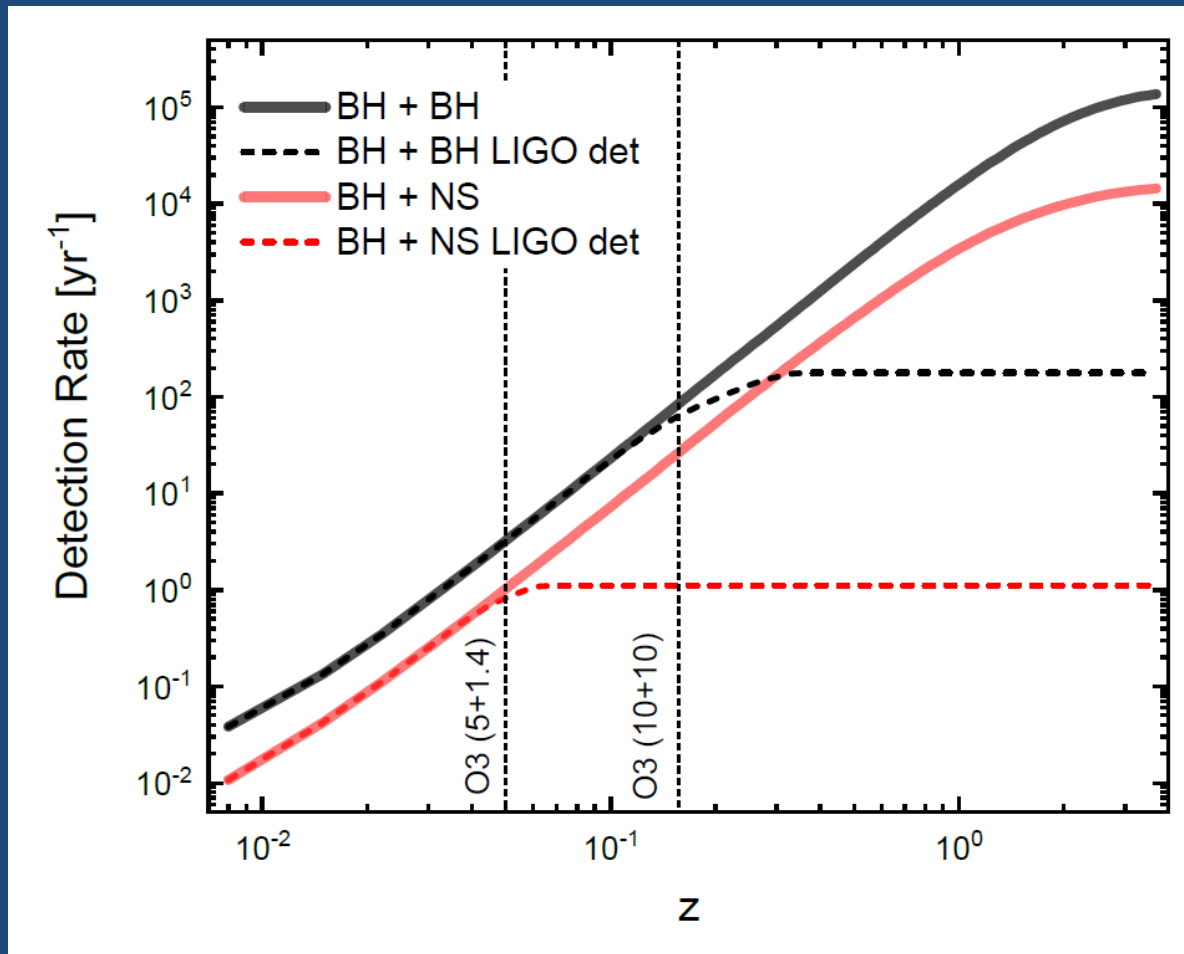


NS+BH upper limit
610 Gpc⁻³ yr⁻¹

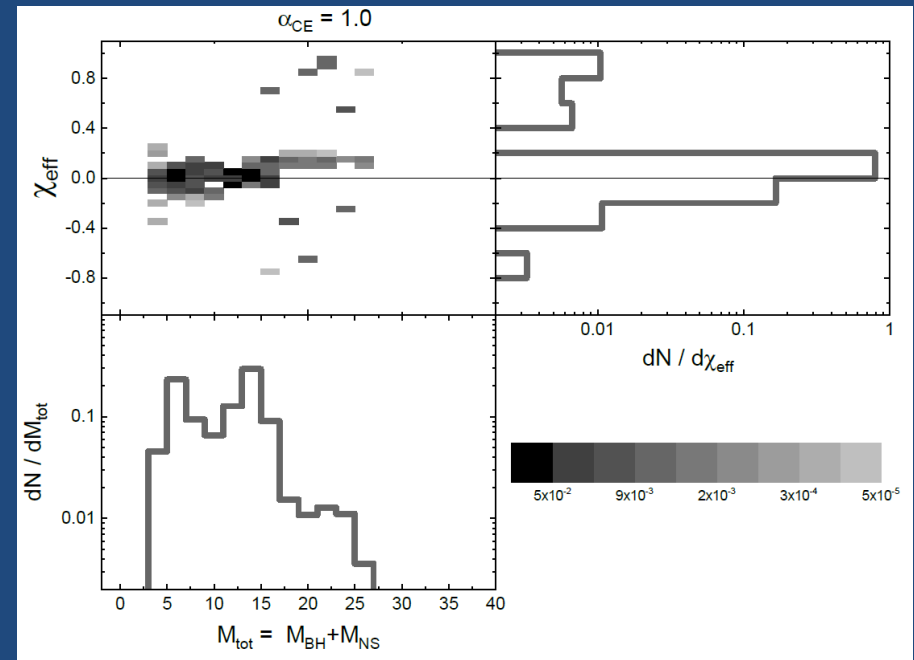
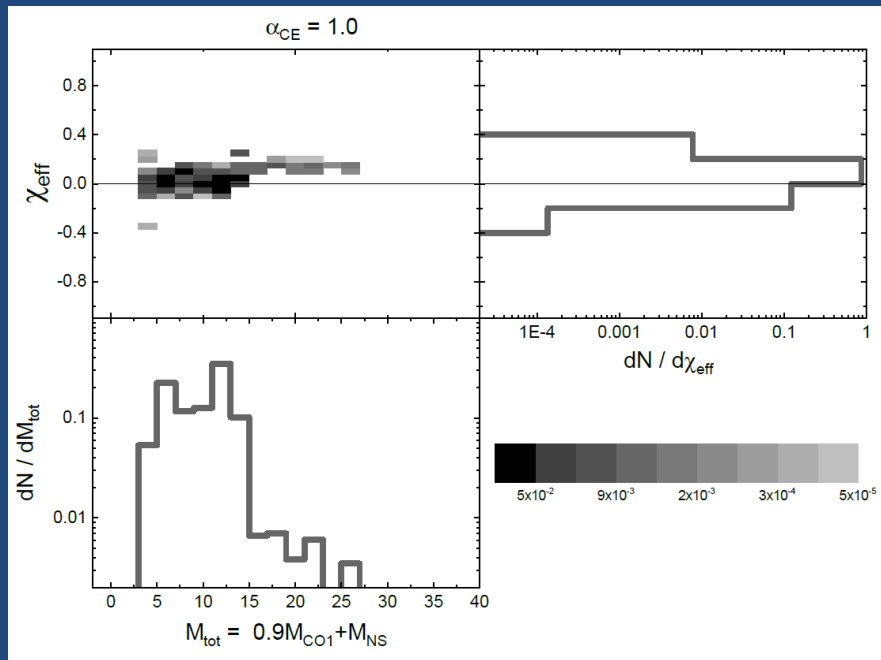
1811.12907

PK+ 2019, Physics-Uspekhi (2019,
No 11, in press); 1907.04218

Detection rate BH+BH, BH+NS



Effective spin/total mass distribution for BH+NS coalescing binaries

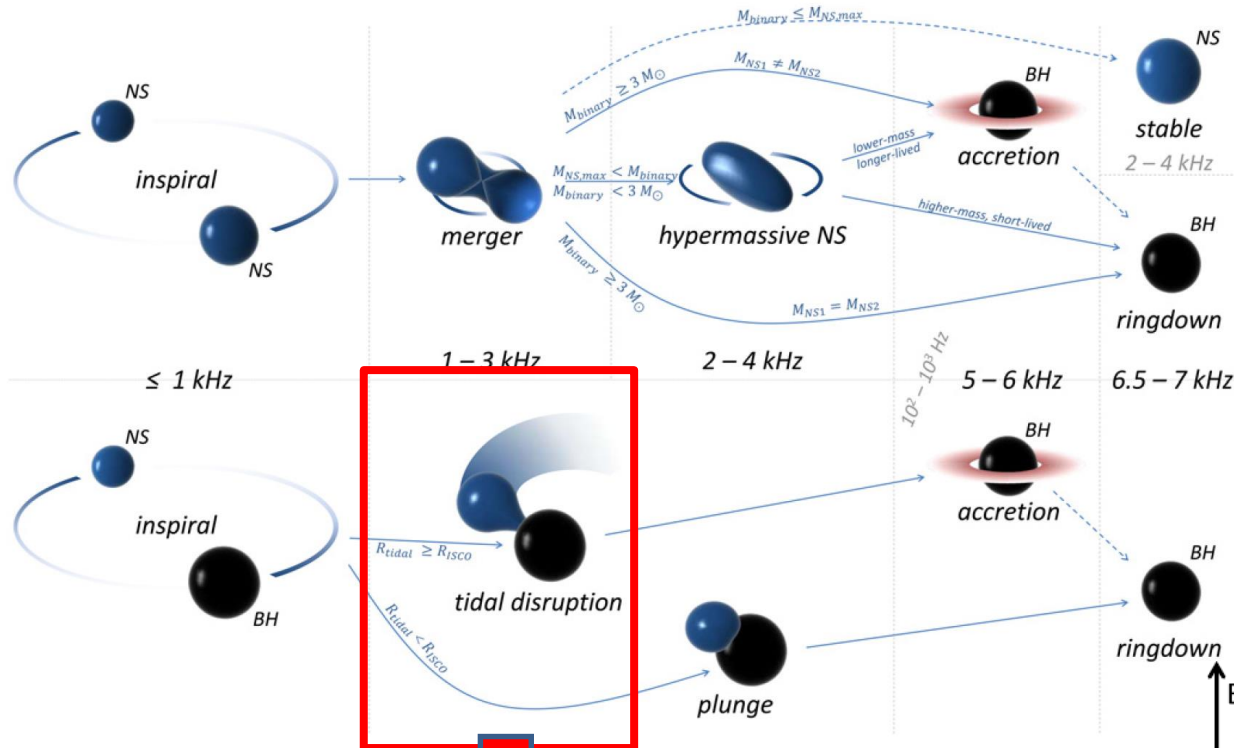


1907.04218

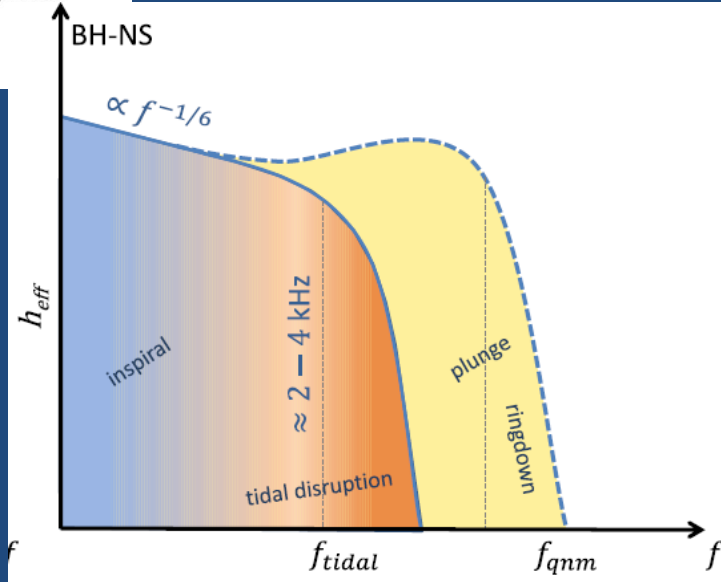
sGRB from NS+BH

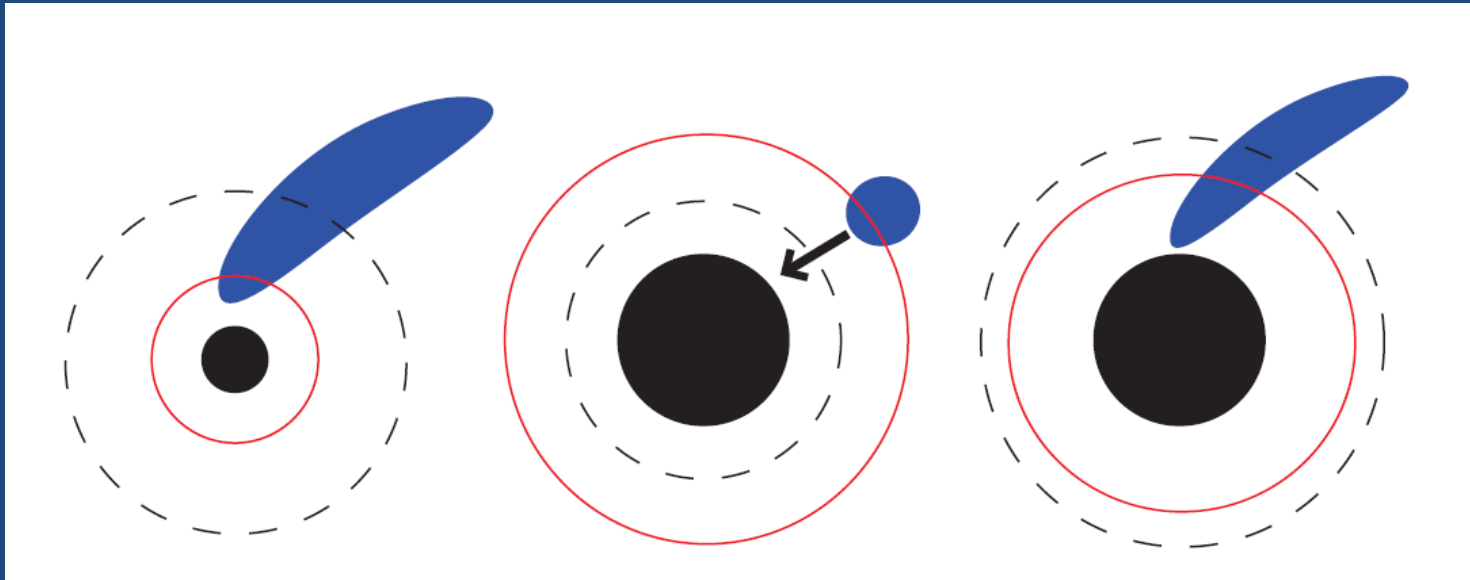
Class. Quantum Grav. **30** (2013) 123001

Topical Review



sGRB-favorable





$$R_{\text{tid}} > R_{\text{ISCO}}$$

$$R_{\text{tid}} < R_{\text{ISCO}}$$

$$R_{\text{tid}} \sim R_{\text{ISCO}}$$

Kyotoku+'11

Mass shedding and tidal disruption

$$r_{\text{ISCO}}/M = 3 + Z_2 \mp \sqrt{(3 - Z_1)(3 + Z_1 + 2Z_2)}$$

$$Z_1 \equiv 1 + (1 - \chi_1^2)^{1/3}$$

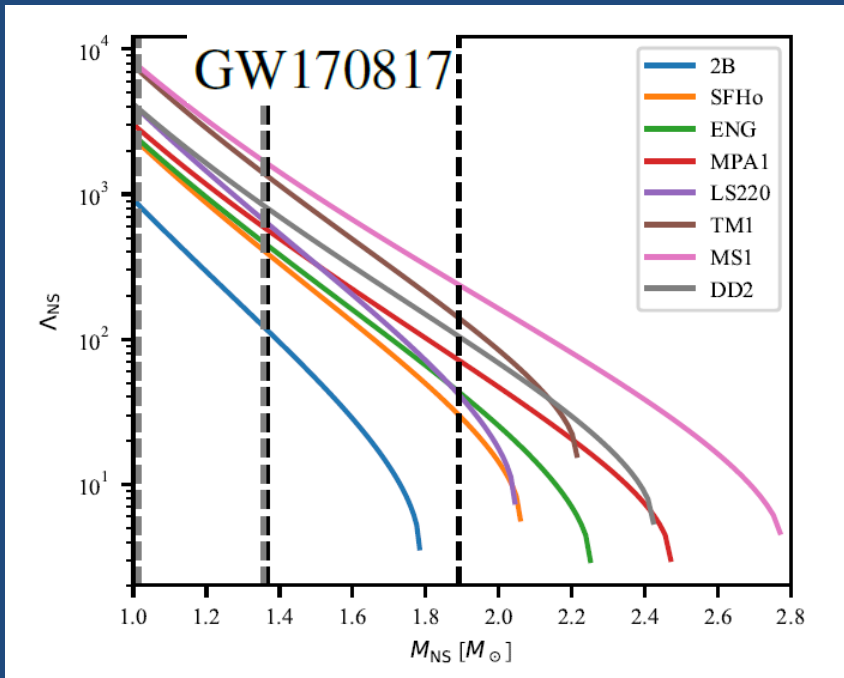
$$\times \left[(1 + \chi_1)^{1/3} + (1 - \chi_1)^{1/3} \right]$$

$$Z_2 \equiv \sqrt{3\chi_1^2 + Z_1^2}$$

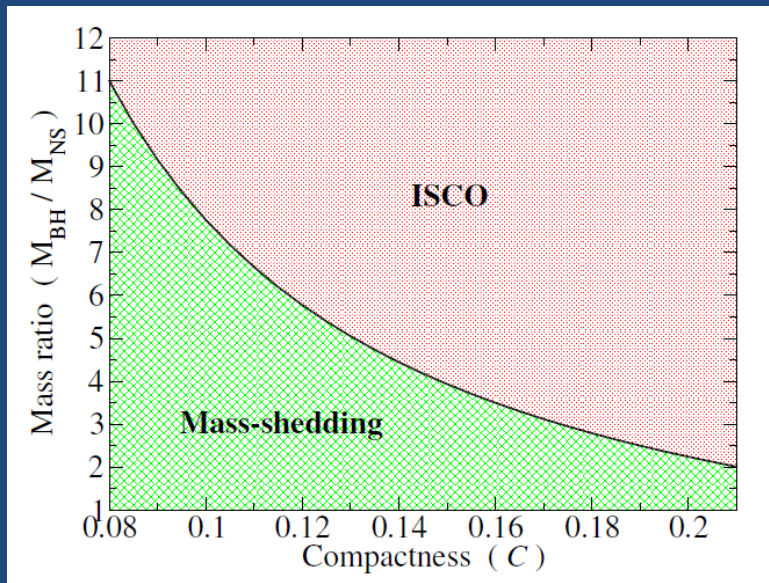
- $R_{\text{tid}} \sim R_{\text{ns}} (M_{\text{bh}}/M_{\text{ns}})^{1/3}$
Mass shedding if

$$R_{\text{tid}} > R_{\text{ISCO}}$$

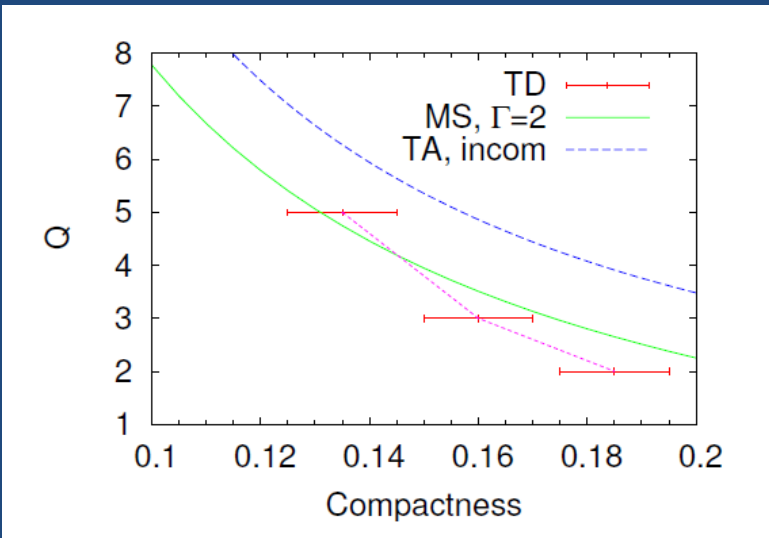
- Depends on NS compactness $C = M_{\text{ns}}/R_{\text{ns}}$ (EOS)
- **Tidal parameter**
 $\Lambda = 2k_2/(3C^5)$
- Depends on the BH spin



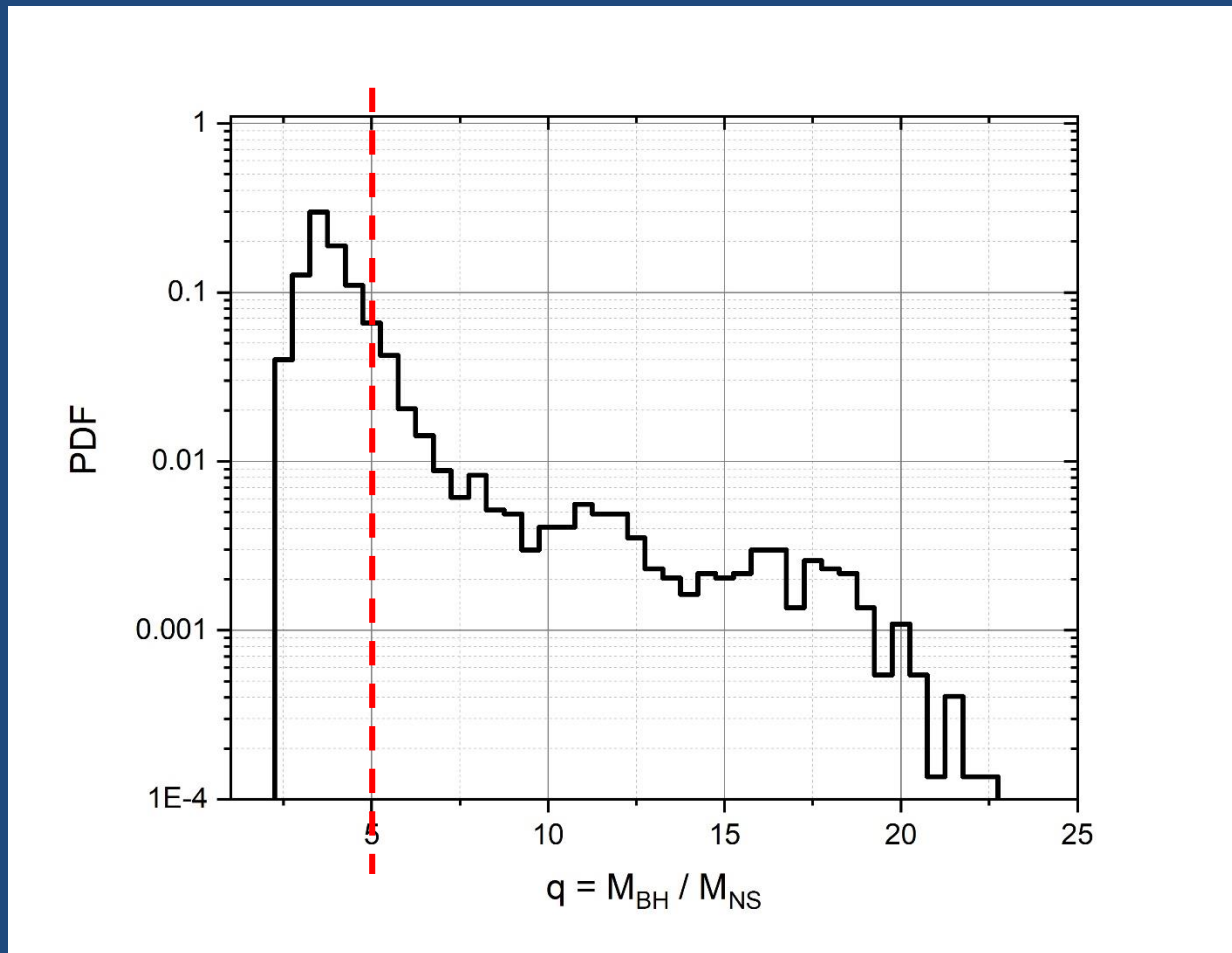
Mass ratio for tidal disruption



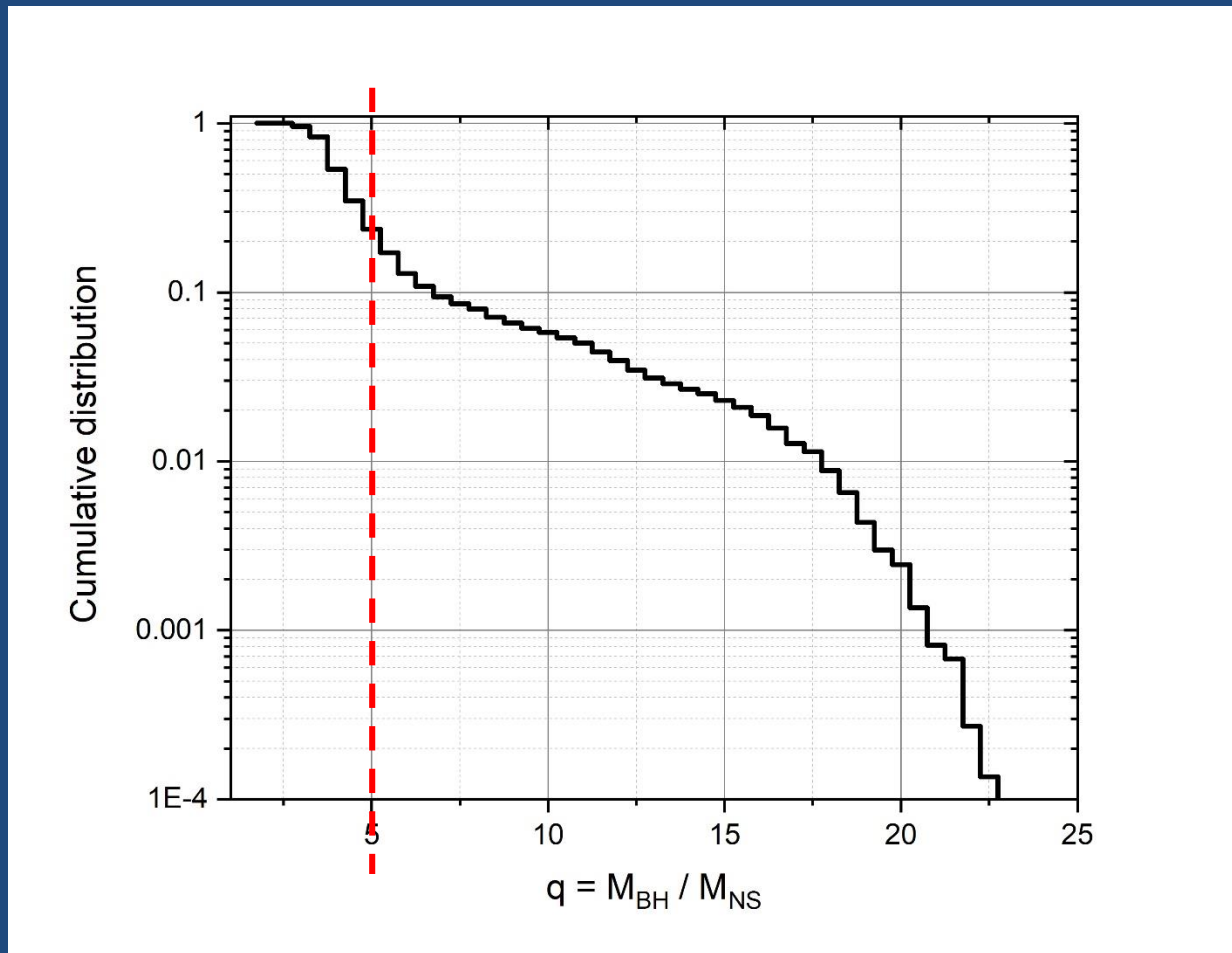
- Approximately the same as for the mass shedding
- For $a=0$: $q_{\text{cr}} \sim 3-5$
- For $a \rightarrow 1$: $q_{\text{cr}} \sim 5-7$



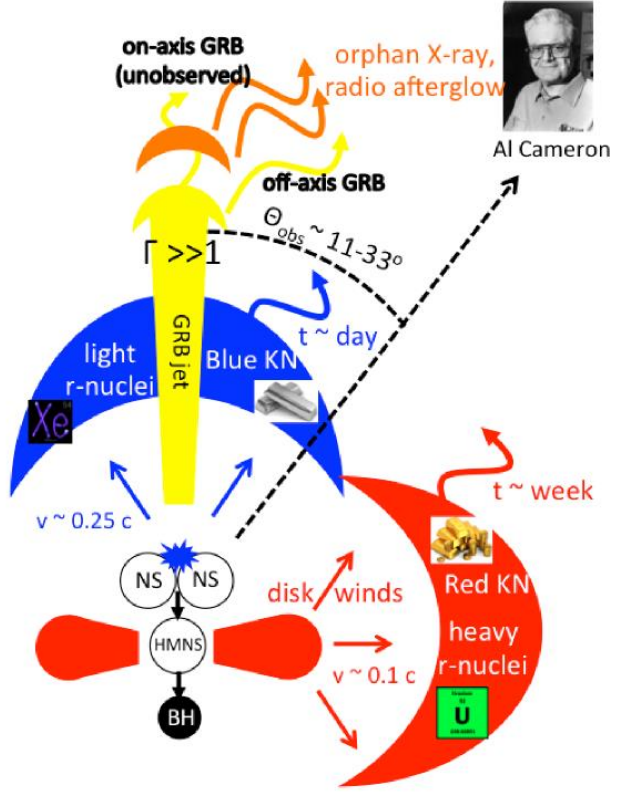
Coalescing BH/NS mass ratio



Coalescing BH/NS mass ratio



- A substantial (>50%) fraction of NS+BH mergings is expected to occur in the tidal disruption regime favoring disk formation around a rotating BH



Mass ejection

- ‘Dynamical’ (merger) + ‘viscous’ (disc)

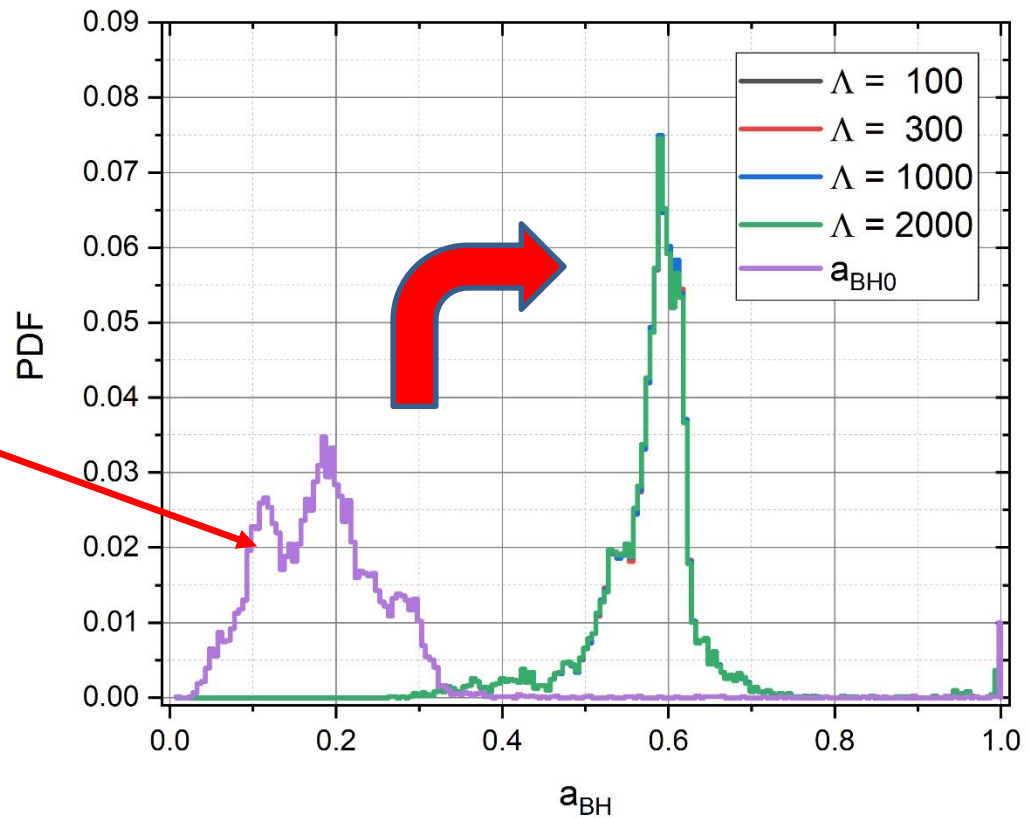
$$Y_e = \frac{n_e}{n_p + n_n} = \frac{n_p}{n_p + n_n}$$

Type of binary	Remnant	$M_{\text{ej,dyn}}$	$M_{\text{ej,vis}}$	$Y_{e,\text{dyn}}$	$Y_{e,\text{vis}}$	$\langle v_{\text{ej}} \rangle$
Low- m BNS	SMNS	$O(10^{-3})$	$O(10^{-2})$	0.05–0.5	0.3–0.5	0.15
Mid- m BNS (stiff EOS)	HMNS	$O(10^{-3})$	$O(10^{-2})$	0.05–0.5	0.2–0.5	0.15
Mid- m BNS (soft EOS)	HMNS	$\sim 10^{-2}$	$O(10^{-2})$	0.05–0.5	0.2–0.5	0.20
High- m BNS ($q \sim 1$)	BH	$< 10^{-3}$	$< 10^{-3}$	—	—	—
High- m BNS ($q \ll 1$)	BH	$O(10^{-3})$	$\lesssim 10^{-2}$	0.05–0.1	0.05–0.3	0.30
BH-NS	BH	0–0.1	0–0.1	0.05–0.1	0.05–0.3	0.30

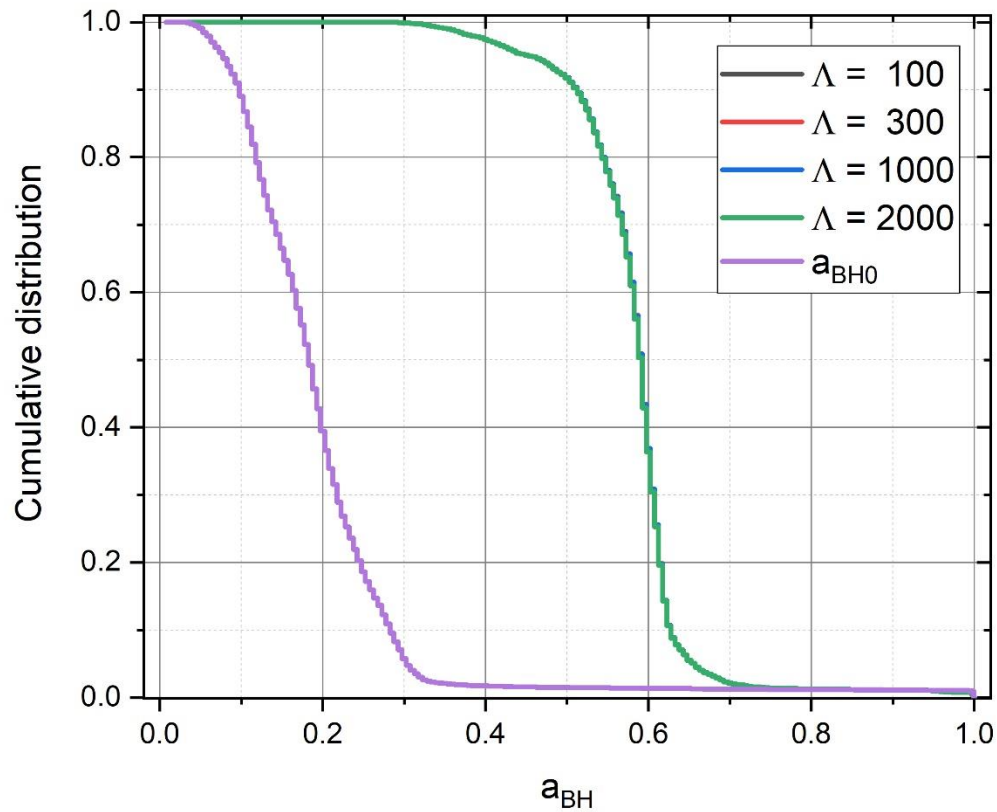
- **Mass ejection depends on the total mass M before the coalescence, binary mass ratio, component spins and tidal deformation (EOS)**
- Final BH mass and spin, emitted GW energy
→ from numerical relativity simulations (Jimenez-Forteza+'18)
- Account for NS EOS → from NR simulations of BH+NS mergings (Zappa+'19)

Final spin of BH

Initial BH spin from population synthesis (PK+'19)

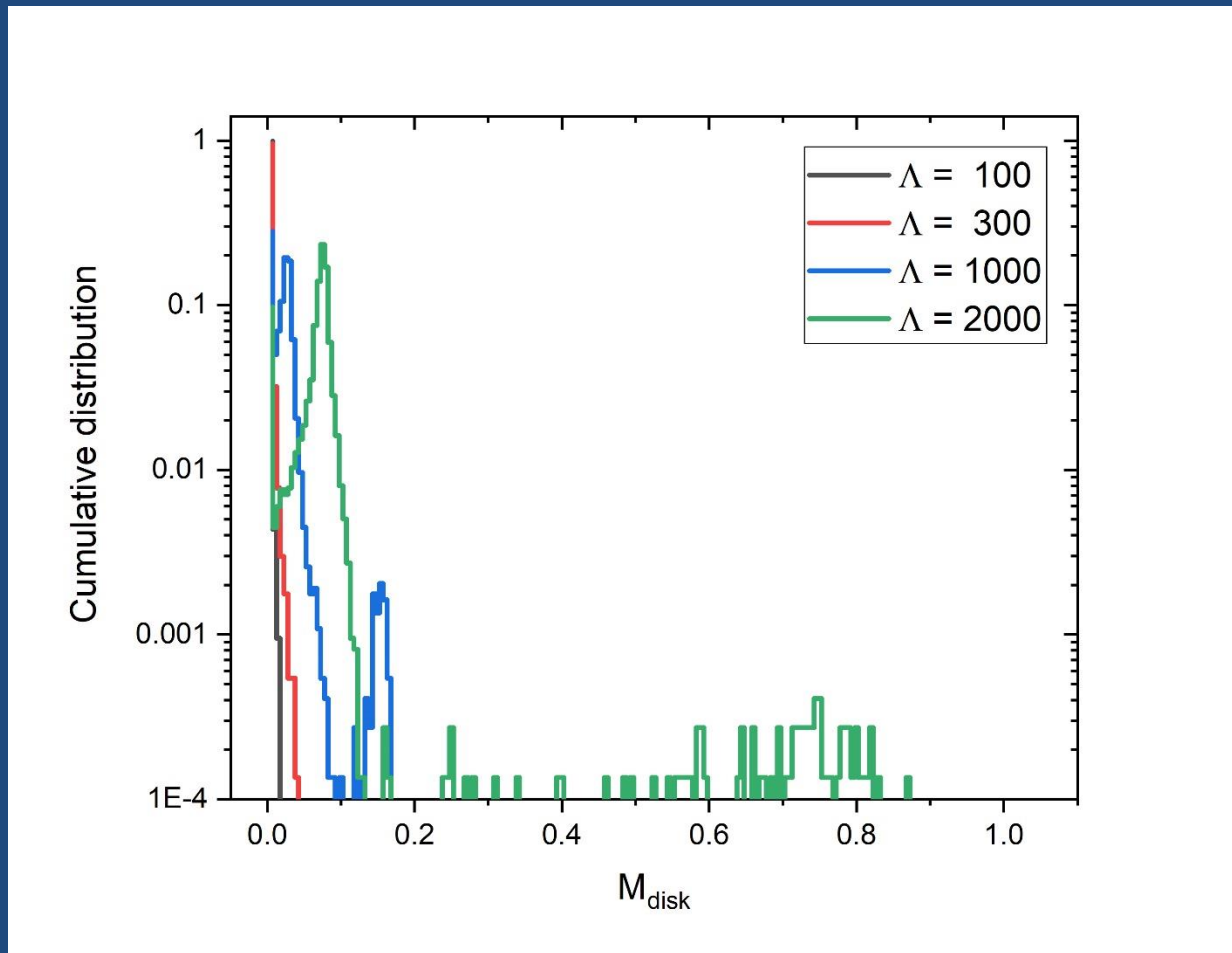


Final spin of BH

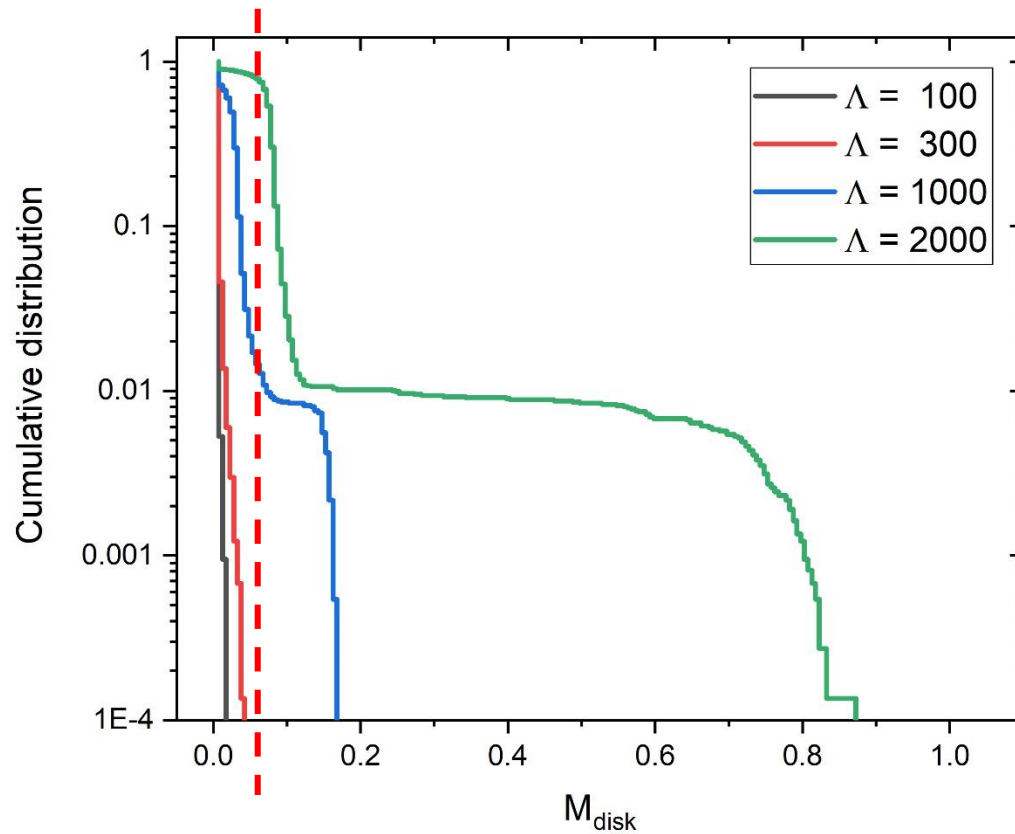


- Final BH spin is almost insensitive to uncertain NS EOS!

Residual disk mass from tidally disrupted NS ($q < 5$)



Residual disk mass



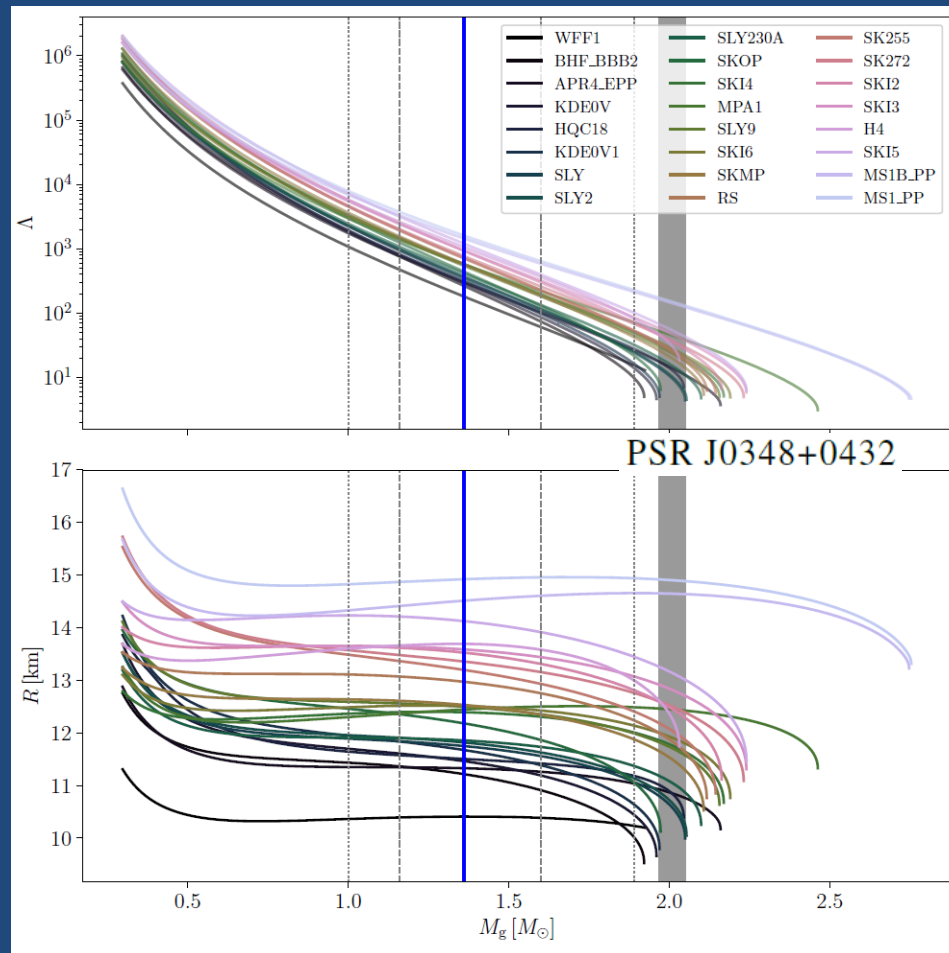
- M_{disk} is determined by the mass ratio
- **Strongly depends on NS EOS!**
- Only large deformations (hard EOS) with **$\Lambda > 1000$** can give rise to interesting disk masses

Constraints from

GW170817:

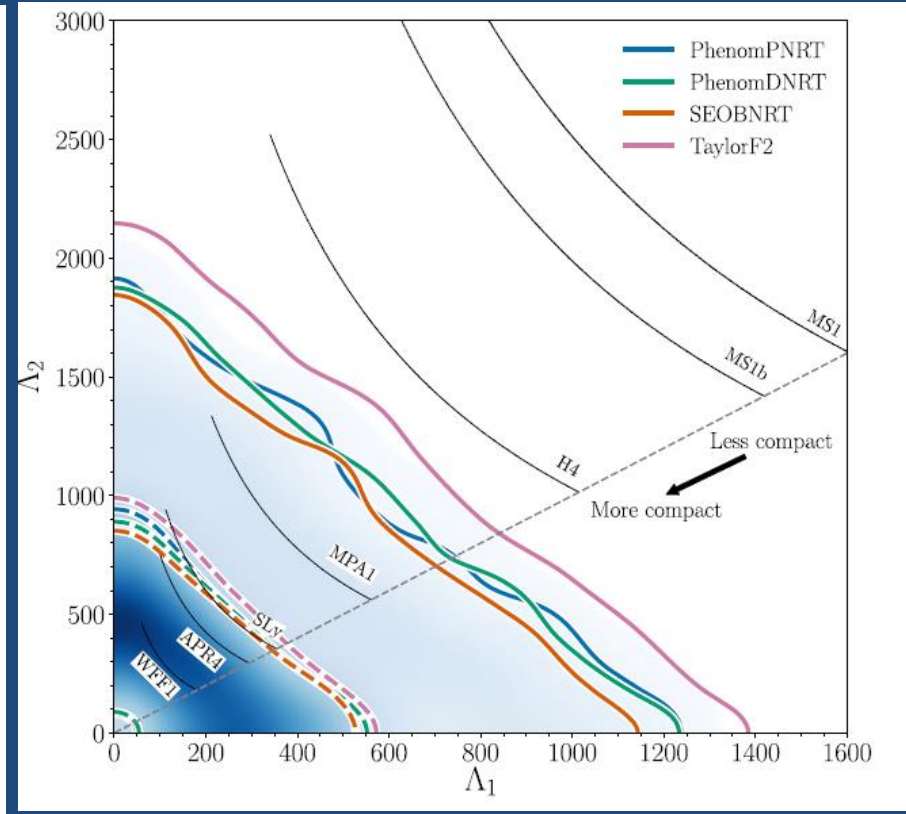
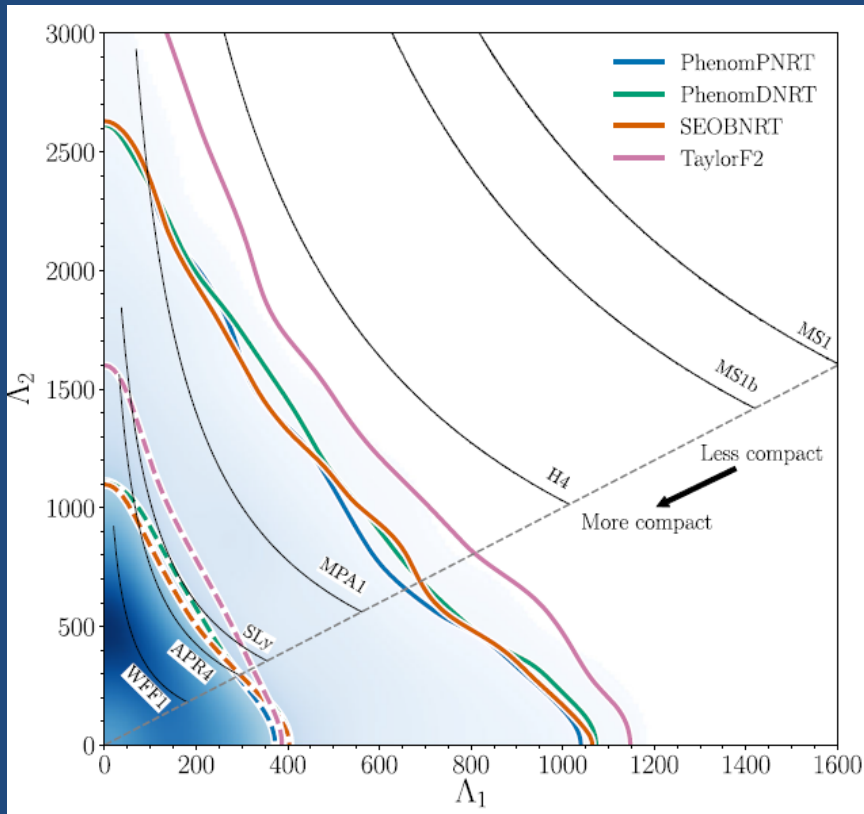
$$\sim 200 < \Lambda < \sim 1600$$

NS equation of state constraints



1908.01012

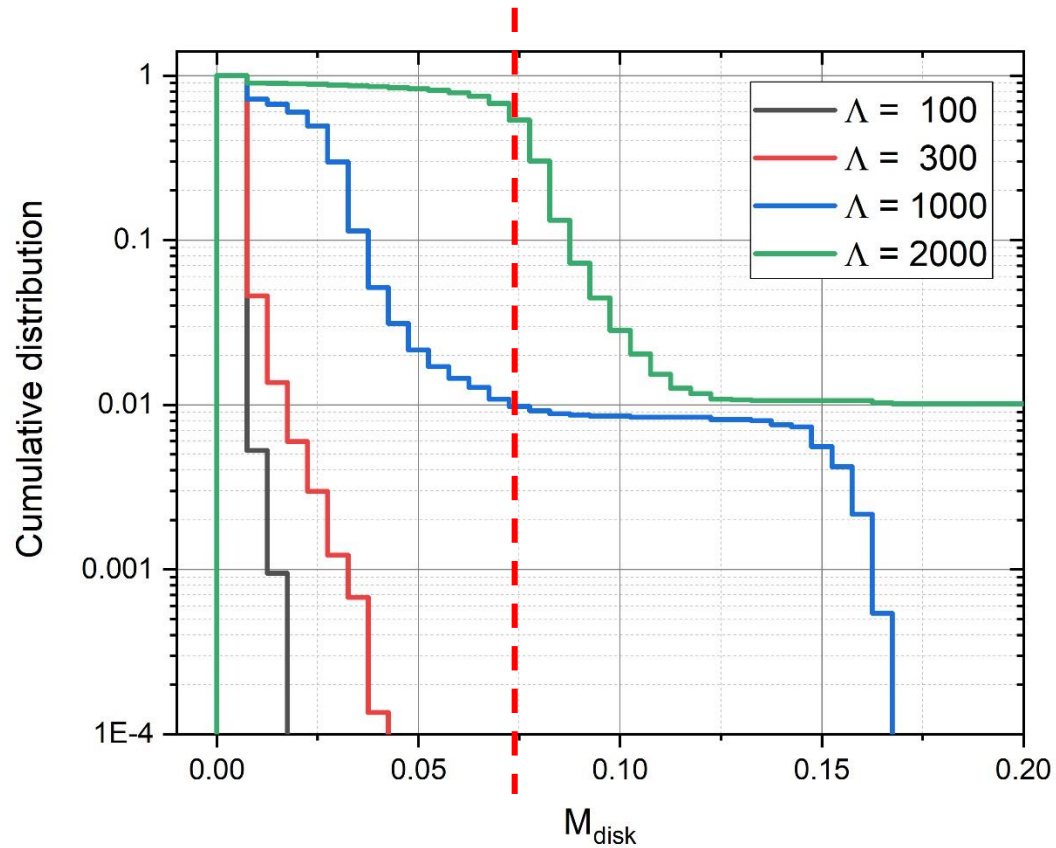
Tidal deformation



$$\Lambda = (2/3)k_2 [(c^2/G)(R/m)]^5$$

1805.11579

Residual disk mass



BZ jet

$$L_{BZ} \sim B_d^2 M^2 \Omega_H^2 f(\Omega_H)$$

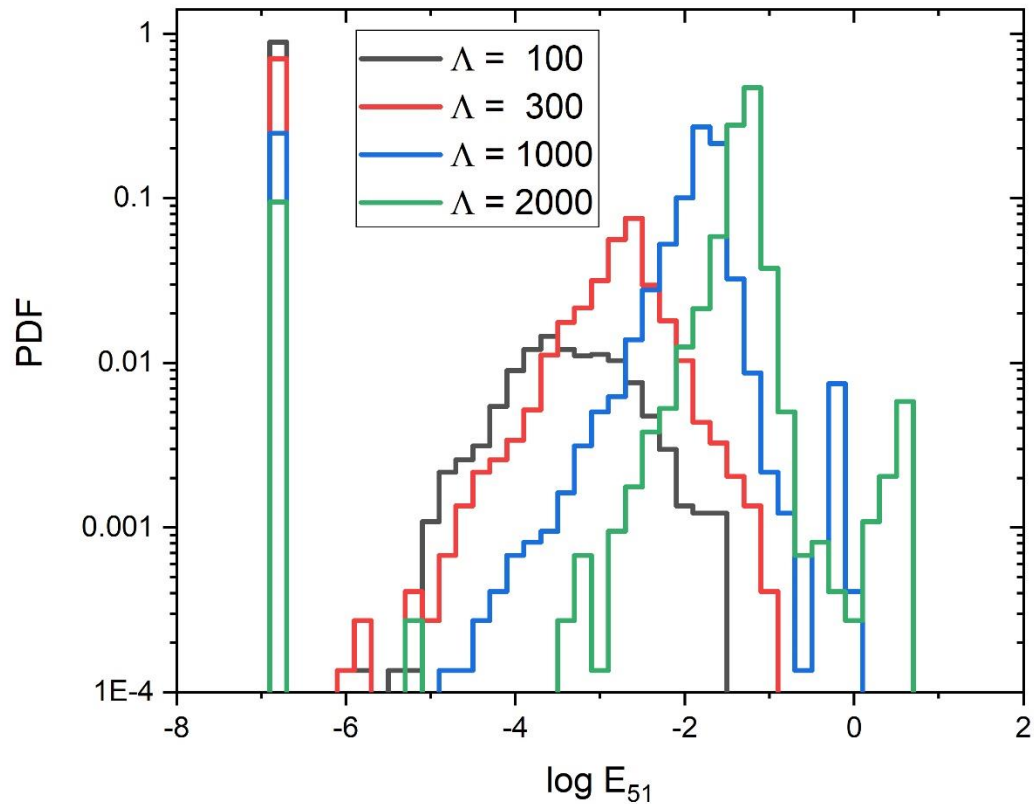
$$B^2 \sim \dot{M}_{accr} / M^2$$

$$\dot{M}_{accr} \sim M_{disk} / t_{accr}$$

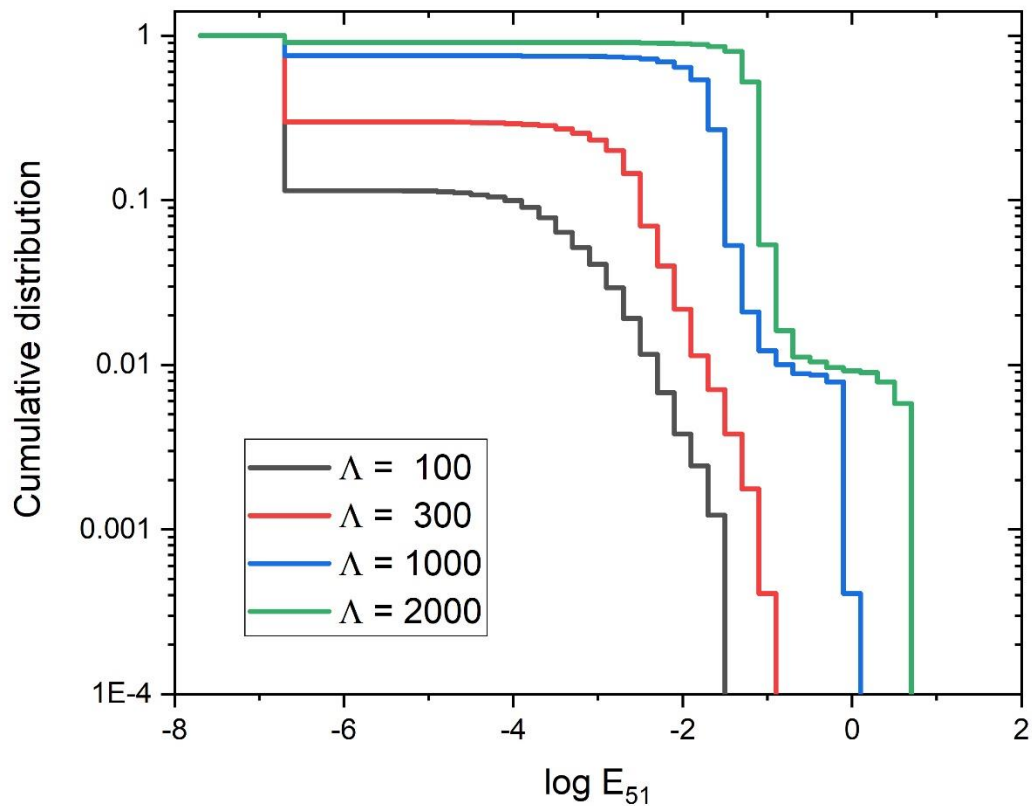
$$E_{BZ} = \varepsilon M_{disk} \Omega_H^2 f(\Omega_H)$$

$$\varepsilon = 0.015 \quad (\text{Barberis} + '19)$$

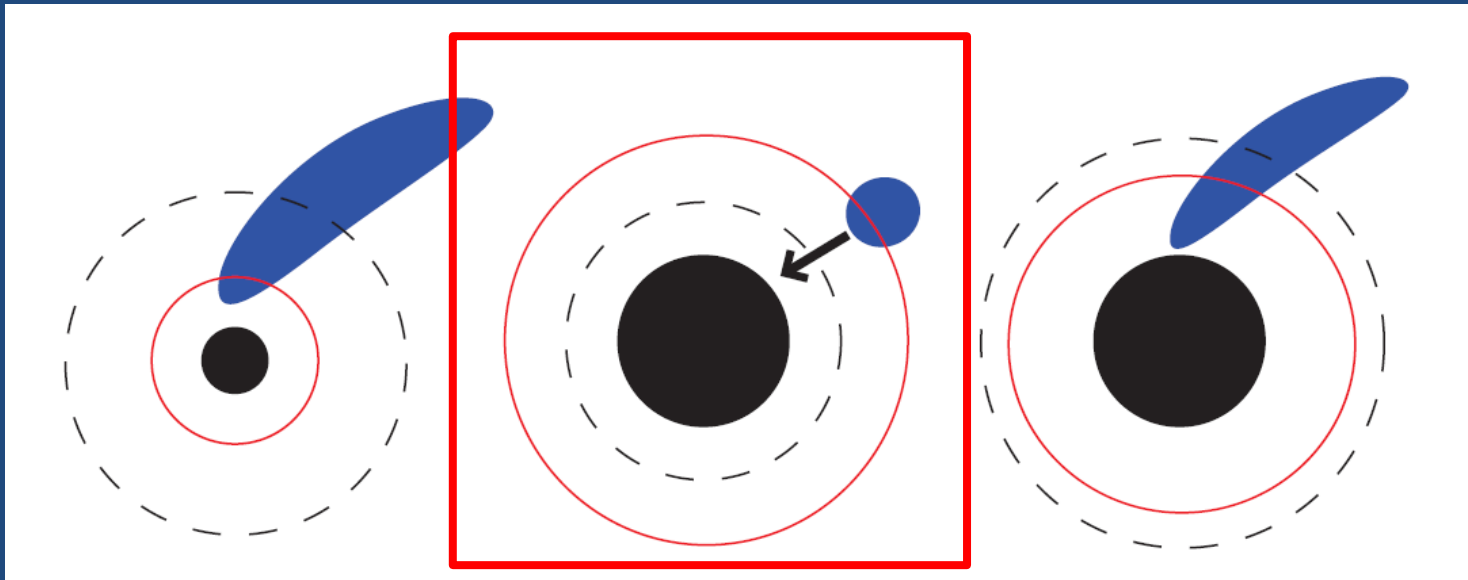
BZ jet kinetic energy



BZ jet kinetic energy



NS plunging into BH ($q > 5$)



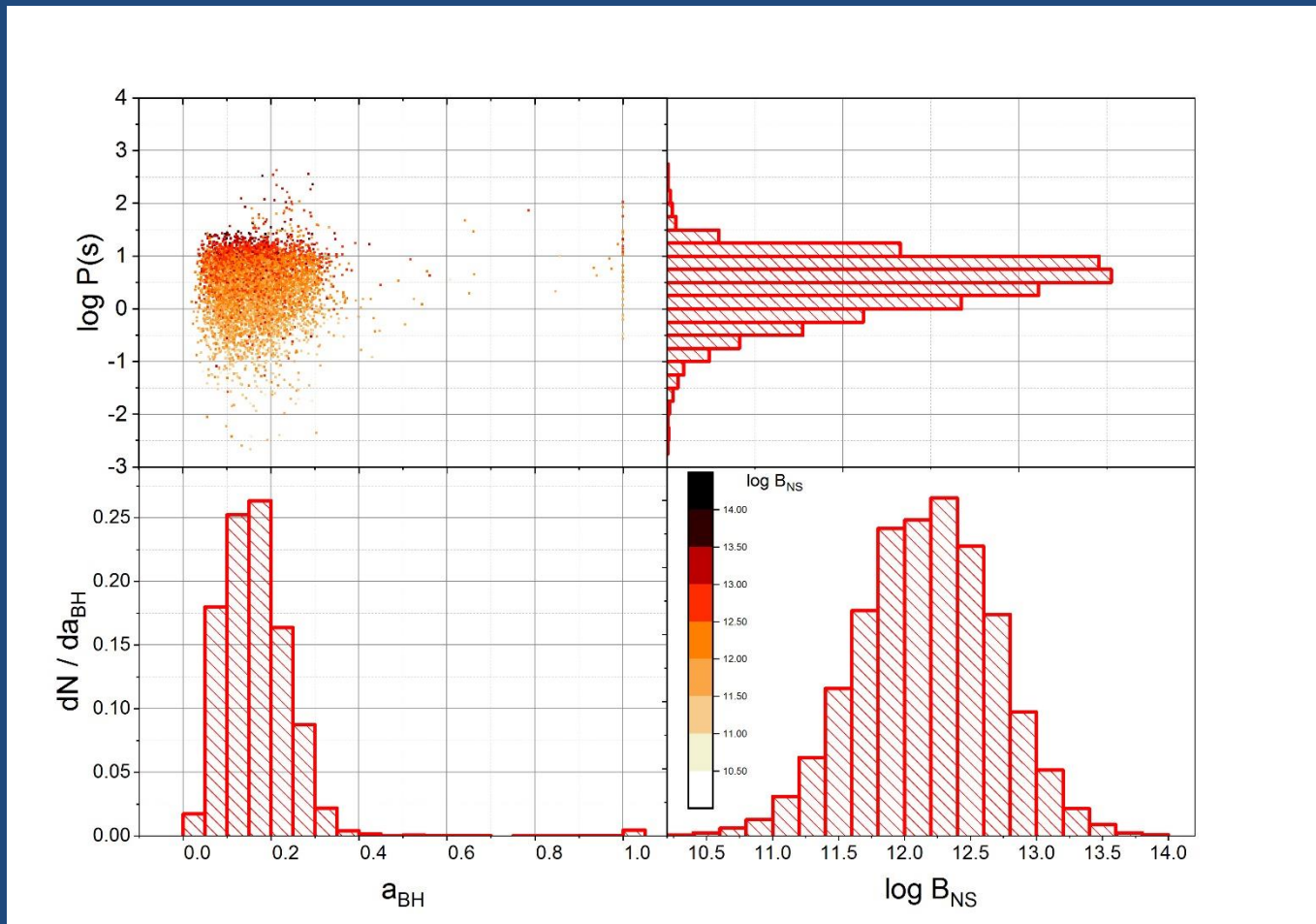
NS plunging into BH ($q > 5$)

- A rotating BH in a magnetic field can acquire electric charge (Wald 1974)

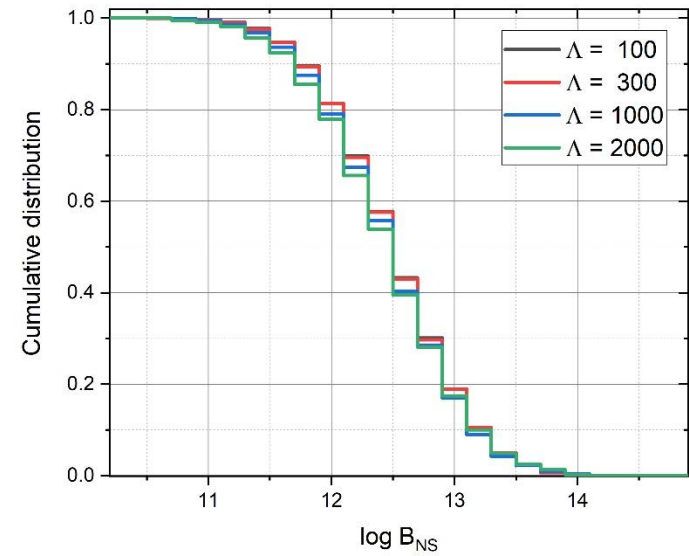
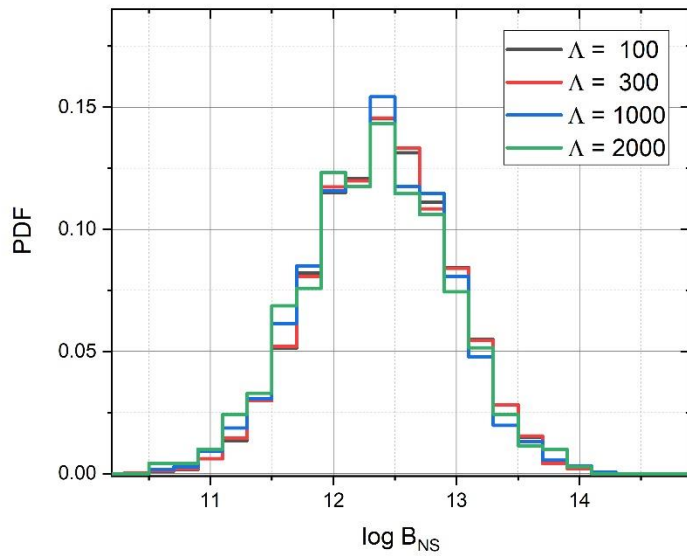
$$Q_{\text{W,max}} \simeq \frac{2G}{c^3} J \times B_{\text{S,NS}} = \frac{2G^2}{c^4} a M^2 B_{\text{S,NS}}$$
$$= 4.4 \times 10^{24} a \left(\frac{M}{10M_{\odot}} \right)^2 \frac{B_{\text{S,NS}}}{10^{12}\text{G}} \text{ e.s.u.,}$$

- There can be EM emission associated with charged BH (Levin+'18, Shu-Quing Zhong+'19...)

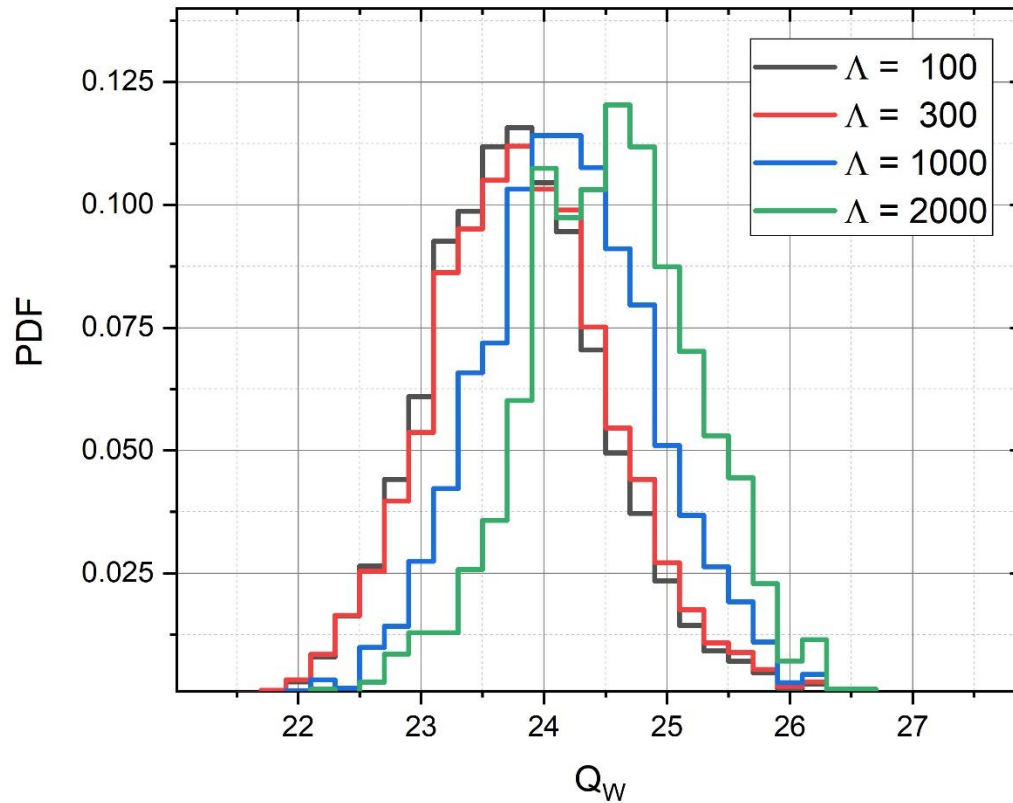
NS rotation and magnetic field before the coalescence



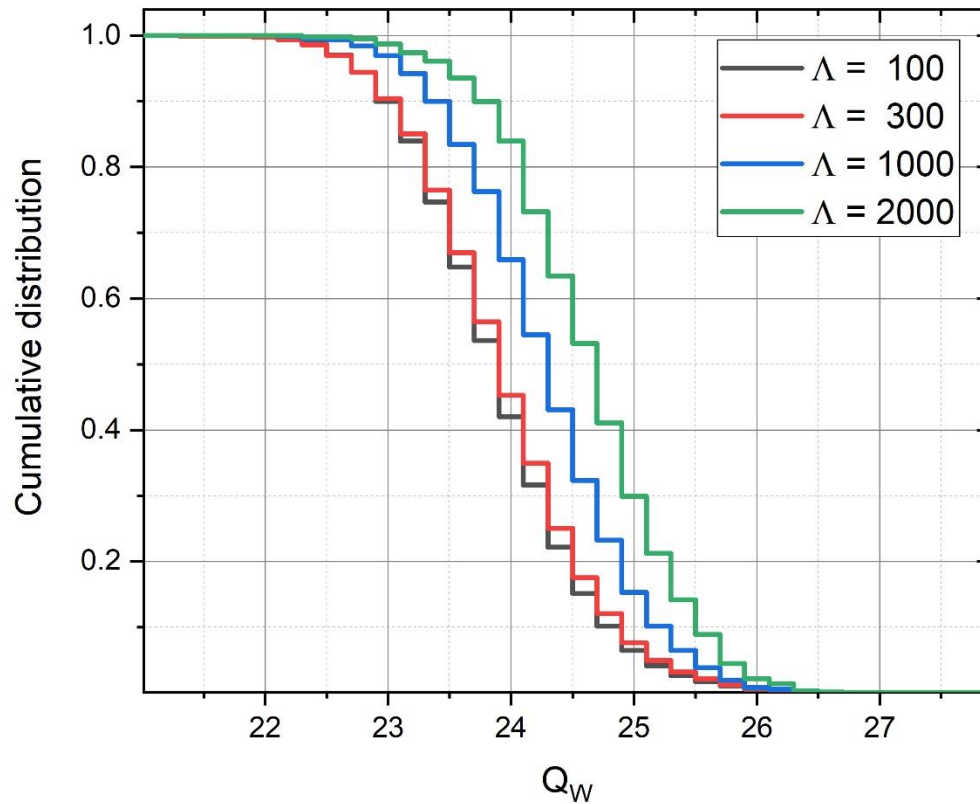
NS magnetic field for plunged NS



Wald BH charge



Wald BH charge



Conclusions

- NS+BH rate is one order of magnitude smaller than BH+BH rate, **~ a few is expected within LVC O3 detection horizon** (as of 13/09/19, 2 out of 25 detections)
- Disk formation from NS tidal disruption mostly depends on (uncertain) NS EOS
- **1-10 % of NS+BH coalescences with tidally disrupted NS can launch relativistic BZ jets and produce short (likely subluminal) GRBs**
- More exotic (but less secure) mechanisms of EM radiation from high-mass-ratio NS+BH plunges are not excluded